

Analytical Soft SUSY Spectrum in Supersymmetric Models in Light of $S_4 \times Z_n$ Flavor, Symmetric SUSY SO(10) Theory.

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ABSTRACT: The heavy right-handed neutrinos in supersymmetric models can act as the source of lepton flavor violation (LFV). LFV processes like $\mu \rightarrow e\gamma$, $\tau \rightarrow \mu\gamma, \tau \rightarrow e\gamma$ is an effective way to explore newphysics beyond the SM. Among the possible processes, μ decays have the greatest discovery potential in most of the supersymmetric models. Experimental inference of lepton flavor-violating processes within a supersymmetric type-II seesaw framework in the non-universal Higgs model (NUHM) and non-universal Scalar Mass model for Yukawa mixing scenarios in the S_4 theory with an additional discrete symmetry is presented. The numerical analysis includes full 2 loop renormalization group running effects for the the above mentioned Yukawa coupling matrices. The projected discovery reach of LFV experiments (MEG-II) is mentioned and those regions in mSUGRA, NUHM, NUSM models that have already been excluded by the LHC searches or that which is probed by MEG experiments is specified here. The results presented in this work can influence experimental challenges and physics motivations to construct various BSM theories and sensitivity to test these theories at next run of HE/HL LHC is also considered.

KEYWORDS: Charged Lepton Flavor Violation, mSUGRA, NUHM, NUSM models.

I. INTRODUCTION

Supersymmetry (SUSY) remains a potent and powerful model for physics Beyond the Standard Model, despite the fact that there is insufficiency of signals from the Large Hadron Collider (LHC) [1]. SUSY solves the naturalness problem i.e, the big hierarchy problem [2], introduces a stable dark matter candidate, neutralino[3] and explains neutrino masses when supplemented with right-handed (RH) Majorana neutrinos [4]. Hitherto, the Minimal Supersymmetric Standard Model (MSSM) [5] has been confronted and yet it has sustained itself into following experimental tests: 1. Gauge coupling unification of three fundamental forces [6], 2. Discovery of the top quark within 100-200 GeV for a successful radiative electroweak symmetry breaking (REWSB) [7], 3. discovery of the Higgs boson $m_h \sim 125$ GeV, within the narrow range of MSSM allowed values [8]. Neutrino oscillations, proved by experiments, requires one to go to physics beyond standard model. These neutrino oscillations, and hence mixings, are also surmised to induce lepton flavor violations in the charged leptonic sector. Theoretically, such cLFV processes could be instigated in different theories with BSM particles such as SUSY GUT[9], SUSY See Saw [10]-[12], LHC Higgs Model [13] and models with extra dimension [14]. In this work cLFV decay $\mu \rightarrow e$, getting contributions from neutrino oscillations and mixings is considered.

The current sensitivity of precision charged lepton flavor violation experiments, cLFV measurements already explores certain regions of SUSY parameter space, mostly the constrained MSSM (cMSSM or mSUGRA) [15] with light scalar masses and right-handed neutrinos (RHN). The MEG Collaboration announced an upper bound for the branching fraction of the process $\mu \rightarrow e$: $BR(\mu \rightarrow e) < 4.2 \times 10^{-13}$ at 90 percent C.L. [16]. The expected sensitivity of the MEG-II experiment is 6×10^{-14} for three years of data taking [17].

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Nevertheless precision experiments cannot be speculated of as a replacement for LHC, they can be compatible. Constructive results for sparticle searches at the LHC in company with LFV search results would constrain many BSM theories that are presently consistent with observations.

Numerous processes implicating lepton flavor violating decays could be feasible such as $\mu \rightarrow e$, $\tau \rightarrow \mu$ or $\tau \rightarrow e$ transitions. Enhancement for $\mu \rightarrow e$ decay at next phase of MEG experiment is expected to reach $BR(\mu \rightarrow e + \gamma) \leq 6 \times 10^{-14}$ [18],[19]. In this work, the decay $\mu \rightarrow e + \gamma$ is considered, as this is top foremostly constrained by experiments. Such experimental hunt, and abstract studies on lepton flavor violation (LFV) can help us compel the new physics or BSM theories, that could be contemporary just above the electroweak scale, or within the jurisdiction of next run of LHC. It is significant, that in the next run of LHC, i.e High Luminosity LHC (HE/HL-LHC), the center of mass energies are conventional to go to 27 TeV [20]. The flavor physics programme at the HL-LHC comprises many different probes- the weak decays of μ , τ leptons and the Higgs boson in which the experiments can search for signs of new physics. It will be exciting to see the full potential of the HE-LHC to serve as a facility for precision new physics for decades to come.

It is needless to say that SUSY GUTs gives intensification to tiny neutrino masses via see saw mechanisms in which noteworthy benefaction to lepton flavor violation could come from flavor violations amidst heavy sleptons. The outcome of lepton flavor violation could become outstanding due to radiative corrections to Dirac Neutrino Yukawa Couplings (DNY), which might become apparent if the see saw scale is slightly lower than the GUT scale [11], [21]-[25]. Alike studies in contrastive see saw mechanisms have been carried out in [11],[21]-[26]. In [11], similar studies

were done in scenario when neutrino masses and mixings appear attributable to type I See Saw mechanism of SUSY SO(10) theory, where the Dirac neutrino Yukawa couplings were of the kind- $Y_v = Y_u$ and $Y_v = Y_u^{\text{diag}} U_{\text{PMNS}}$, where

$Y_u = V_{\text{CKM}} Y_u^{\text{diag}} V_{\text{CKM}}^\dagger$. In this work Similar studies are done in type II See Saw scenario in $S_4 \times Z_n$ flavor symmetric model. Lepton Flavor Violation in SUSY type II seesaw [27] models have also been considered untimely in [23]-[25]. In this work examination on LFV decay ($\mu \rightarrow e$) using type II see saw mechanism in $S_4 \times Z_n$ flavor symmetric SUSY SO(10) theories [28] is carried out, and hence the reactivity to try out the surveillance of sparticles at HE/HL run of LHC [20], in mSUGRA, NUHM, and NUSM [39] models is also discussed. Such studies in Non Universal Gaugino Mass models were done earlier in [30, 31]. It is noteworthy that $S_4 \times Z_n$ flavor symmetric SUSY SO(10) theory gives correct fit to observed neutrino oscillations and mixings and provides specific mass textures for the quarks and leptons with only a small number of parameters and also predicts quark lepton mass relations and mixing angles in both the quark and the lepton sector. Specially, the model leads to tri-bi-maximal form for the PMNS matrix in the leading order with corrections to this imminent from charged lepton fields. In [11, 31], similar studies were done employing type I see saw formula, applying older value of $BR(\mu \rightarrow e)$ [32]. The structure of Dirac neutrino Yukawa couplings from [28] is used here in this work, for $\tan\beta = 10 - 60$, and $M_{\text{GUT}} = 2 \times 10^{16} \text{ GeV}$. The value of Higgs mass as measured at LHC [20] and improved precision values of reactor mixing angle θ_{13} as measured at Daya Bay, Reno [33] have been utilised in this work. Few studies on LFV in SO(10) GUTs have also been conferred in [34],[35]. The minimal supergravity model (mSUGRA) is a well driven and induced model [36], in which Minimal Supersymmetric Standard Model (MSSM) [37] can be inserted. In mSUGRA, SUSY is broken in the hidden sector, and is proclaimed to the visible sector MSSM fields by dint of gravitational interactions. Formation of gaugino masses [38] in mSUGRA (N=1 supergravity) embraces two scales - spontaneous SUGRA breaking scale in the hidden sector over the singlet chiral superfield and the another one is GUT breaking scale via the non singlet chiral superfield [36]. In postulate these two scales can be distinct. But in a minimalistic perspective, they are usually assumed to be similar [36]. This conducts to a common mass m_0 for all the scalars, a common mass $M_{1/2}$ for all the gauginos and a common trilinear SUSY breaking term A_0 at the GUT scale, $M_{\text{GUT}} \approx 2 \times 10^{16} \text{ GeV}$.

Now, the universality of sfermion masses, assumed in mSUGRA, NUHM models is analysed here. SO(10) symmetric soft terms approximately mean boundary conditions close to NUHM. In the scheme of SO(10) theories, all the matter

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fields, and the right handed neutrino, are there in the same 16-dimensional representation, and consequently, all the matter fields will have identical mass at the high scale. Nevertheless, the higgs fields can have dissimilar mass, as they do not survive in the same representation as the matter field. Therefore, the boundary terms for the SO(10) theory are compatible with NUHM and mSUGRA (in mSUGRA, all the higgs will be in equivalent representation). Diverging from NUHM boundary conditions, will generally gesture a deviation from SO(10) boundary conditions. If the hidden sector has representations which are not singlets under SO(10), one can expect non-trivial gaugino mass boundary conditions. So, to sum up, both NUHM and NUGM are boundary conditions which are a result of assuming SO(10) symmetric boundary conditions at the GUT scale in two disparate ways. Moreover, as can be perceived from Figs. 4 presented in section IV, low energy flavor phenomenology is not much attacked by these distinct boundary conditions at high scales.

In this work a SUGRA model with non universal scalar masses [39] where the first two generations of scalar masses and the third generation of sleptons are extremely massive is also probed. This model implores the flavor changing neutral current (FCNC) issue by subscribing very large masses for the first two generations of squarks and sleptons. But the need of radiative breaking of electroweak symmetry REWSB prohibits the scalar masses from being too hugely massive. This situation is eluded by allowing third generation squark masses and the Higgs scalar mass parameters to be of small scale. [39]. This smallness also set out to keep the naturalness problem within sovereignty. It is conventional that SUSY can be broken by soft terms of kind $-A_0, m_0, M_{1/2}$, where A_0 is the universal trilinear coupling, m_0 is the universal scalar mass, and $M_{1/2}$ is the universal gaugino mass. Stern universality amongst Higgs and matter fields of mSUGRA models can be reduced or weakened in NUHM (Non Universal Higgs Mass [40] models. As manifested in the results in Sec.IV in mSUGRA, the spectrum region of $M_{1/2}$ and m_0 is found to settle in the direction of heavy side, as enabled by MEG constraints on $BR(\mu \rightarrow e)$, but in NUHM, lighter spectra is feasible (owing to partial cancellations in flavor violating terms). So it enthrals one to study LFV decay $\mu \rightarrow e$ in NUSM (Non Universal Scalar Mass Models) [31, 39]. As manifested in the results in sec.IV it is found that in NUSM model the gaugino masses are very massive, so as to enable very large scalar masses. In the light of the third generation squark masses and the Higgs scalar mass parameters being small, the fine tuning problem of naturalness does not get deteriorated. So to have Higgs mass around 125.9 GeV, the first two generation of squark and slepton masses as well as third generation of slepton masses habitats around 12000 GeV-16000 GeV. It is found that in NUSM model, current value of the branching fraction of the process $\mu \rightarrow e$: $BR(\mu \rightarrow e) < 4.2 \times 10^{-13}$ at 90 percent C.L. allows $\tan \beta$ to lie in the region 5–45 and in order to have Higgs mass around 125.9 GeV, NUSM model grants $\tan \beta$ to find its appropriate value in the region 20–30 whereas in NUHM model it is seen that $\tan \beta$ occupies itself around 7–13 for Higgs mass around 125.9 GeV.

Therefore it is conceived that indication of LFV could be tested at high luminosity HE/HL run of LHC, if SUSY particles are spotted around TeV range. It is eminent that, no SUSY partner of SM has been noticed yet at LHC, and thistipst to a high scale SUSY theory. The LHC has rigorous limits on sparticles, which could entail a tuning of EW symmetry at a few percent level [41]-[46]. Thus some substitute to low scale SUSY theories have been recommended. Few of them are – minisplit SUSY [47] and maximally natural SUSY [48]. In minisplit SUSY the scalar sparticles are massive than the sfermions (gauginos and higgsinos), so that sfermions could be spotted at HE-LHC. Scalar sparticles could be available everywhere in the range $(10-10^5)$ TeV. In maximally natural SUSY, the 4D theories rises from 5D SUSY theory, with Scherk-Schwarz SUSY breaking at a Kaluza-Klein scale $\sim 1/R$ of several TeV[48]. Some prospects of LFV in those theories have been explored in [49]. The paper has been organised as follows. In section II, connections of LFV with type II See Saw mechanism in $S_4 \times Z_n$ flavor symmetric SUSY SO(10) theory is discussed. In section III, the values of various parameters used in this analysis has been presented. The software SuSeFLAV[50] is used to determine $BR(\mu \rightarrow e)$. Section IV manifests itself in results and their analysis. Section V outlines the work.

II. LEPTON FLAVOR VIOLATION IN $S_4 \times Z_N$ FLAVOR SYMMETRIC SUSY SO(10) THEORY.

A. Lepton Flavor Violation and observables.

In SUSY theory, non-diagonal mass matrix elements in the slepton mass matrix are the initiator of Lepton Flavor Violation processes.

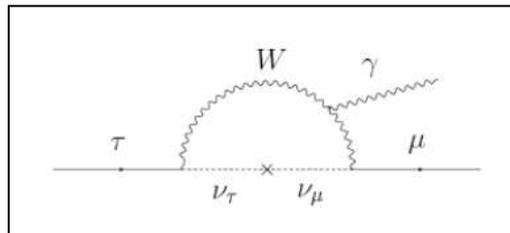


Fig.1. Examples of Feynmann Diagrams contributing to $\tau \rightarrow \mu + \gamma$ processes in SUSY models.

The SUSY SO(10) theory obviously embraces the seesaw mechanism. The existence of heavy RH neutrinos at an intervening scale give on to the running and originate flavor violating entries in the left-handed slepton mass matrix at the weak scale [11]. The lepton flavour violating entries in the SO(10) SUSY GUT framework can be perceived in terms of the low energy parameters. In the mass insertion (MI) method with leading log approximation, the branching fractions for different LFV processes $l_i \rightarrow l_j + \gamma$ can be resembled as:

$$BR(l_i \rightarrow l_j + \gamma) \approx \alpha^3 \frac{|(m_{(L)}^2)(LL)_{ij}|^2}{G_F^2 m_{SUSY}^8} \tan^2 \beta BR(l_i \rightarrow l_j \nu_i \bar{\nu}_j) \tag{1}$$

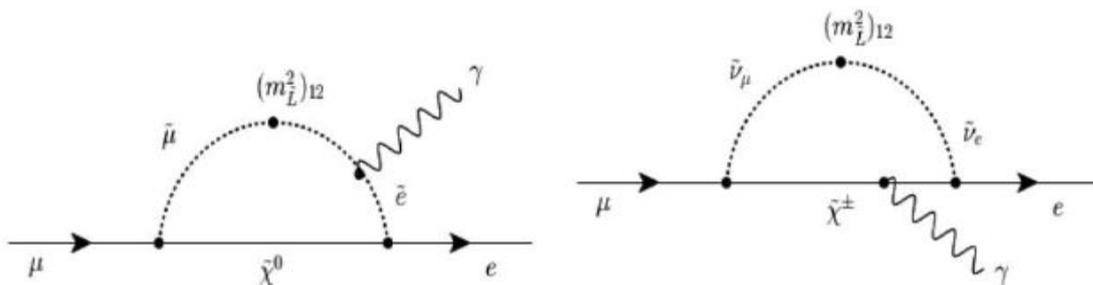


Fig. 2.Examples of Feynmann Diagrams contributing to $\mu \rightarrow e + \gamma$ processes in SUSY models.

where M_{SUSY} is mas scale of the SUSY particles, α is the fine structure constant and GF is the Fermi constant. The 6×6 slepton mass matrix.

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Vol. 9, Issue 3, March 2020

$$m_{(L)}^2 = \begin{pmatrix} m_{(L)}^2(LL)_{ij} & m_{(L)}^2(LR)_{ij} \\ m_{(L)}^2(RL)_{ij} & m_{(L)}^2(RR)_{ij} \end{pmatrix} \quad (2)$$

where LL, RL, LR, RR are 3×3 entries established on the chirality tag of sfermions. The leading log entries approximation in mSUGRA is defined as [51] is

$$(m_{\tilde{L}}^2)_{i \neq j} = \frac{-3m_0^2 + A_0^2}{8\pi^2} \sum_k (f_\nu^*)_{ik} (f_\nu)_{jk} \log \left(\frac{M_X}{M_{Rk}} \right) \quad (3)$$

where M_X is the GUT scale, M_{Rk} is the k th heavy RH majorana neutrino scale, m_0 and A_0 are universal soft mass and trilinear terms at the high GUT scale. f_ν are the Dirac neutrino Yukawa (DNY) couplings. The flavour violation is specified in terms of the quantity $\delta_{ij} = \Delta_{ij}/m_1^2 m_i^2$ is the geometric mean of the slepton squared masses [52], and $\Delta_{i \neq j}$ are flavour non diagonal entries of the slepton mass matrix instigated at the weak scale due to RG evolution. The mass insertions are labelled into the LL/LR/RL/RR kinds [53], based on the chirality of the SM fermions. The fermion masses are formed by renormalisable Yukawa couplings of the $10 \oplus 126 \oplus 120$ portrayal of scalars of SO(10) GUTs. The Dirac neutrino Yukawa couplings f_ν at the high GUT scale in $S_4 \times Z_n$ flavor symmetric SUSY SO(10) theory is used in this work from [28].

$$f_\nu = \frac{1}{v \sin \beta} M_D \quad (4)$$

Here M_D is the Dirac neutrino mass matrix. The off-diagonal flavor violating entries at the weak scale in eq. (3) are then absolutely decided by using f_ν from eq. (4). The δ s from the RGEs are calculated, using the leading log approximation. The soft SUSY masses are flavour universal at the input scale, so, off diagonal entries in the LL sector are developed by running of right handed neutrinos in the renormalisation group equation loops. For using the leading log expression (eq. (3)) one requires the mass of the heaviest right handed neutrino, which is here, $\sim 10^{16}$ GeV. The induced off-diagonal leading log entries pertinent to LFV decay $l_i \rightarrow l_j + \gamma$ are of the magnitude of (setting down A_0 to 0)[31].

Table I: Governing values of δ_{ij} that enter eq. (5,6,7) for in $S_4 \times Z_n$ flavor symmetric SUSY SO(10) theory.

LFV contributions For $S_4 \times Z_n$ flavor symmetric SUSY SO(10) theory case	
δ_{12}	0.6672×10^{-4}
δ_{23}	1.5634×10^{-4}
δ_{31}	0.7377×10^{43}

In NUHM models, the expression $(-3m_0^2 + A_0^2)$ of mSUGRA models in eq.(3) is put back by the leading log approximation for the slepton mass matrix element that induces the process $\mu \rightarrow e + \gamma$, as $(-2m_0^2 + A_0^2 + m_{H_u}^2)$.

International Journal of Innovative Research in Science, Engineering and Technology

(A High Impact Factor, Monthly, Peer Reviewed Journal)

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Vol. 9, Issue 3, March 2020

m_{Hu} is the soft mass terms of the up type Higgs at the high scale. At the GUT scale, the NUHM case, $m_{Hu} = m_{Hd}$. Besides, due to a relative sign difference between the universal soft mass terms for the matter fields and the Higgs mass terms at the GUT scale, cancellations occurs for $m_{Hu}^2 \approx -2m_0^2$, or increment for $m_{Hu}^2 \geq m_0^2$ in contrast to mSUGRA for the off diagonal flavor violating entries at the weak scale.

$$(\delta_{LL})_{\mu e} = \frac{-3}{8\pi^2} (f_\nu^*)_{13} (f_\nu)_{23} \ln \left(\frac{M_X}{M_{R_3}} \right) \quad (5)$$

$$(\delta_{LL})_{\tau \mu} = \frac{-3}{8\pi^2} (f_\nu^*)_{23} (f_\nu)_{33} \ln \left(\frac{M_X}{M_{R_3}} \right) \quad (6)$$

$$(\delta_{LL})_{\tau e} = \frac{-3}{8\pi^2} (f_\nu^*)_{13} (f_\nu)_{33} \ln \left(\frac{M_X}{M_{R_3}} \right) \quad (7)$$

B. $S_4 \times Z_n$ flavor symmetry

Dirac Neutrino Yukawa couplings from a supersymmetric $SO(10)$ grand unified theory (GUT) of flavor relevant to an S_4 family symmetry is used here in this work. It makes use of the fact that, $SO(10)$ theory combined with type II seesaw mechanism for generating neutrino masses mingled with a simple assertion that the dominant Yukawa coupling matrix (the 10-Higgs coupling to matter) has rank one. The rank one model arises within some reasonable speculation as a constructive field theory from vectorlike 16 dimensional matter fields with masses lying above the GUT scale. S_4 flavon multiplets get vevs in the ground state of the theory. By enlarging the S_4 theory with an additional discrete symmetry Z_n , it has been found that the flavon vacuum field alignments acquire a discrete values of parameters assuming that some of the higher dimensional couplings are small. An observed set of vacuum alignments directs one to an unification of quark-lepton flavor: (i) the lepton mixing matrix that is dominantly tri-bimaximal with small corrections related to quark mixings; (ii) quark lepton mass is related at GUT scale as $m_b = m_\tau$ and $m_\mu = 3m_s$ and (iii) the solar to atmospheric neutrino mass ratio $m_{\text{solar}}/m_{\text{atm}} = \theta_{\text{Cabibbo}}$, which in tally with the experiments. The mass matrix becomes

$$M_\nu = \begin{pmatrix} 0 & c & c \\ c & a & c-a \\ c & c-a & a \end{pmatrix} \quad (8)$$

where $c/a = \sqrt{2} \leq 1$. It is diagonalized by the tri-bi-maximal matrix.

International Journal of Innovative Research in Science, Engineering and Technology

(A High Impact Factor, Monthly, Peer Reviewed Journal)

Visit: www.ijirset.com

Vol. 9, Issue 3, March 2020

$$U_{TB} = \begin{pmatrix} \sqrt{\frac{2}{3}} & \sqrt{\frac{1}{3}} & 0 \\ -\sqrt{\frac{1}{6}} & \sqrt{\frac{1}{3}} & -\sqrt{\frac{1}{2}} \\ -\sqrt{\frac{1}{6}} & \sqrt{\frac{1}{3}} & \sqrt{\frac{1}{2}} \end{pmatrix} \quad (9)$$

This is not the full U_{PMNS} matrix which is going to get small radiative corrections from diagonalization of the charged lepton mass matrix, which generates small reactor angle, θ_{13} along with small θ_{\odot} and θ_{atm} . The neutrino masses are given by $m_{\nu_3} = 2a - c$; $m_{\nu_2} = 2c$ and $m_{\nu_1} = -c$. To fit observations, $c/a = 0.2$, which realises the neutrino masses to be, $m_{\nu_3} = 0.05\text{eV}$, $m_{\nu_2} = 0.01\text{eV}$, and $m_{\nu_1} = 0.005\text{eV}$.

III. BRANCHING RATIO, $BR(\mu \rightarrow e + \gamma)$ IN CMSSM, NUHM, NUSM

In this section the computations on the LFV constraints of $BR(\mu \rightarrow e)$ in $S_4 \times Z_n$ flavor symmetric SUSY $SO(10)$ theory with type II Seesaw mechanism using the NUHM, CMSSM, NUSM like boundary conditions through detailed numerical analysis is presented. The soft parameter space for CMSSM in the following ranges is studied.

$$\tan\beta \in [1, 60]$$

$$\Delta m_H \in 0$$

$$m_0 \in [0, 8] \text{ TeV}$$

$$M_{1/2} \in [0.3, 4.5] \text{ TeV}$$

$$A_0 \in [-3m_0, +3m_0]$$

$$\text{sgn}(\mu) \in \{-, +\} \quad (10)$$

The analytical part is done using the publicly available package SuSeFLAV[50]. LFV for the non universal Higgs model without completely universal soft masses at high scale is scrutinised. Range of examination of various SUSY parameters, used in NUHM are:

International Journal of Innovative Research in Science, Engineering and Technology

(A High Impact Factor, Monthly, Peer Reviewed Journal)

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Vol. 9, Issue 3, March 2020

$$30 \text{ GeV} \leq m_0 \leq 8 \text{ TeV}$$

$$30 \text{ GeV} \leq M_{1/2} \leq 5 \text{ TeV}$$

$$-9.5 \text{ TeV} \leq m_{H_u} \leq +9.5 \text{ TeV}$$

$$-9.5 \text{ TeV} \leq m_{H_d} \leq +9.5 \text{ TeV}$$

$$-24 \text{ TeV} \leq A_0 \leq +24 \text{ TeV} \tag{11}$$

The Δ_{ij}^{LL} owing to non universal Higgs and $m_h \geq 125 \text{ GeV}$ puts a powerful restraint on SUSY parameter space. Due to partial cancellations in the entrance of Δ_{ij}^{LL} in NUHM case, a substantial region of parameter space can be probed by MEG.

Contingent scans for the following range of parameters in NUSM model is performed [39] and the SUSY particle spectrum using the publicly available package SuSeFLAV[50] is created.

$$\tan\beta \in [5, 60]$$

$$m_0 \in [0, 16] \text{ TeV}$$

$$M_{1/2} \in [0, 6] \text{ TeV}$$

$$A_0 \in 0 \text{ TeV}$$

$$m_{H_u} = m_{H_d} \in 0 \text{ TeV} \tag{12}$$

Massive right handed neutrinos used in our calculations are - $M_{R1} = 10^{13} \text{ GeV}$, $M_{R2} = 10^{14} \text{ GeV}$, and $M_{R3} = 10^{16} \text{ GeV}$. The central values of Δm_{sol}^2 , Δm_{atm}^2 and θ_{13} , from the recent global fit of neutrino data [33] is employed. In Table 1 the presiding values of Δ_{ij} that enter eq.(5,6,7) is presented.

International Journal of Innovative Research in Science, Engineering and Technology

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Visit: www.ijirset.com

Vol. 9, Issue 3, March 2020

IV. CALCULATIONS AND DISCUSSION ON RESULTS

In this section, study on the computation of results presented in section 3 is discussed.

A. Complete Universality – CMSSM

At the high scale, the parameters of the CMSSM model described are m_0 , A_0 and unified gaugino mass $M_{1/2}$. Also there are the Higgs potential parameter μ and the ratio of the Higgs VEVs, $\tan\beta$. The overall SUSY mass spectrum is estimated once those parameters are accessible. The updated MEG constraint [16, 17] set together with a big θ_{13} [33] brings forth significant restrictions on SUSY parameter space in CMSSM. As depicted from fig 3a, few part of the parameter space is allowed for $\tan\beta = 5-60$ in CMSSM as constrained by future MEG limit for $BR(\mu \rightarrow e)$ which is 6×10^{-14} . So to conclude it is seen that the parameter space $M_{1/2} \geq 10$ GeV is permitted by present MEG II bounds on $BR(\mu \rightarrow e)$, incidentally future MEG limit prohibits strictly almost whole $M_{1/2}$ space. Similarly as seen from fig 3b, the current MEG limit allows very heavy spectra for m^0 , which favours m^0 to lie between 4.5 TeV to 8 TeV which is indeed massive. The permitted space in fig 4a needs very massive spectra, i.e. $m^0 \geq 6$ TeV and it mostly lies around 8 TeV for the whole spectrum of $M_{1/2}$. From fig 4b, it is found that the current bound for the process $BR(\mu \rightarrow e)$ 4.2×10^{-13} sets very stringent constraint on $\tan\beta$, which find its value from around 2 to 10. $\tan\beta \leq 2$ and $\tan\beta \geq 10$ is ruled out in the CMSSM case. In fig 4c, 4d the light Higgs mass, m_h as a function of m_0 , $M_{1/2}$ in the CMSSM case is computed. The range of Higgs mass as given by the data at LHC, i.e. $123 \text{ GeV} \leq m_h \leq 127 \text{ GeV}$ allows m_0 $6 \text{ TeV} \geq 8 \text{ TeV}$ as restricted by constrained stringent MEG bounds on $BR(\mu \rightarrow e)$. The region $M_{1/2} \geq 3 \text{ TeV}$ is permitted as can be seen from fig 4d. The asymmetry in A_0 GeV can be seen from fig 4g. To have Higgs mass around, $123 \text{ GeV} \leq m_h \leq 127 \text{ GeV}$, the values of $\tan\beta$ from 3 to 6 are mostly allowed. The dependency of A_0 on $\tan\beta$ is shown in fig 4g.

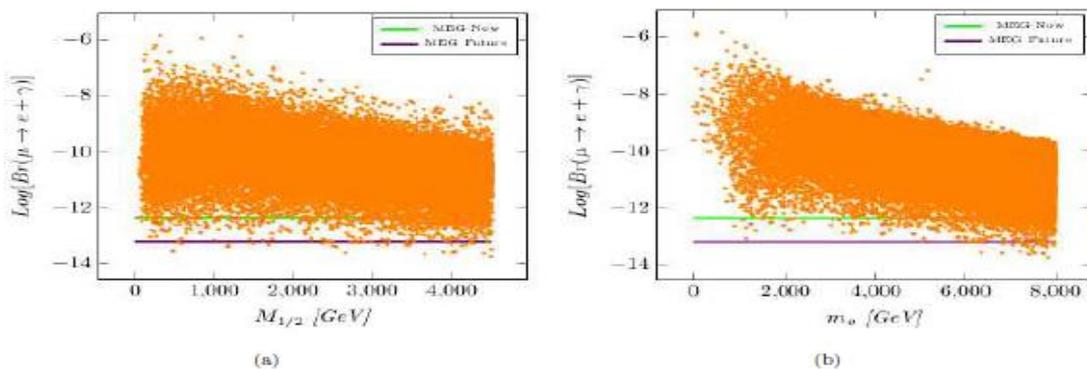


Figure 3: The outcome of the calculations are presented for CMSSM case. In fig 3a, 3b, different horizontal lines depicts the present (MEG 2016) and future MEG constraints for $BR(\mu \rightarrow e + \gamma)$.

B. Non Universal Higgs Model (NUHM)

The cMSSM+RHN model parameter set is investigated in the literature [54–56]. Also, the NUHM1 +RHN models were extensively studied for specific non-universal scenarios with the following GUT-scale mass relations: NUHM1 ($m_{H_u}^2 = m_{H_d}^2 = m_{20}^2$) in [40] The leading log approximation for the slepton mass matrix element that induces the process $\mu \rightarrow e + \gamma$ is:

International Journal of Innovative Research in Science, Engineering and Technology

(A High Impact Factor, Monthly, Peer Reviewed Journal)

Visit: www.ijirset.com

Vol. 9, Issue 3, March 2020

$$(m_L^2)_{i \neq j} = \frac{-2m_o^2 + A_o^2 + m_{H_u}^2}{8\pi^2} \sum_k (f_\nu^*)_{ik} (f_\nu)_{jk} \log \left(\frac{M_X}{M_{R_k}} \right) \quad (13)$$

For low values of $M_{1/2}$, the approximation holds good but for values of $M_{1/2} \approx 1$ TeV, the branching fractions may be up to a factor of 10 dissimilar than the evaluation from RGE running. If the hierarchy of dirac neutrino yukawa couplings are alike to that of the up-type in the Standard Model, where the third generation predominantly dominates, then the largest donation is from the $k = 3$ terms in the summation. Calculation of $\text{Log}[\text{BR}(\mu \rightarrow e + \gamma)]$, in the NUHM case leads to a good approximation for a light/pre-LHC SUSY mass spectrum due to the cancellations between the soft SUSY parameters in eq.(13). Alongside, the consequences of the results accessed in NUHM case is discussed. In fig. 5a m_0 [GeV] Vs $\text{Log}[\text{BR}(\mu \rightarrow e + \gamma)]$, and the fig. 5b in the right panel depicts $M_{1/2}$ [GeV] Vs $\text{Log}[\text{BR}(\mu \rightarrow e + \gamma)]$. Various horizontal lines in fig. 5a, 5b correspond to present and future bounds on $\text{BR}(\mu \rightarrow e + \gamma)$. It is see from the figs. 5a and 5b that due to the partial cancellations between the NUHM parameters m_0 , A_0 and m_{H_u} , almost all of the NUHM parameter space is going to be probed by the stringent MEG bounds for a good approximation of a light/pre-LHC SUSY mass spectrum.

International Journal of Innovative Research in Science, Engineering and Technology

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Vol. 9, Issue 3, March 2020

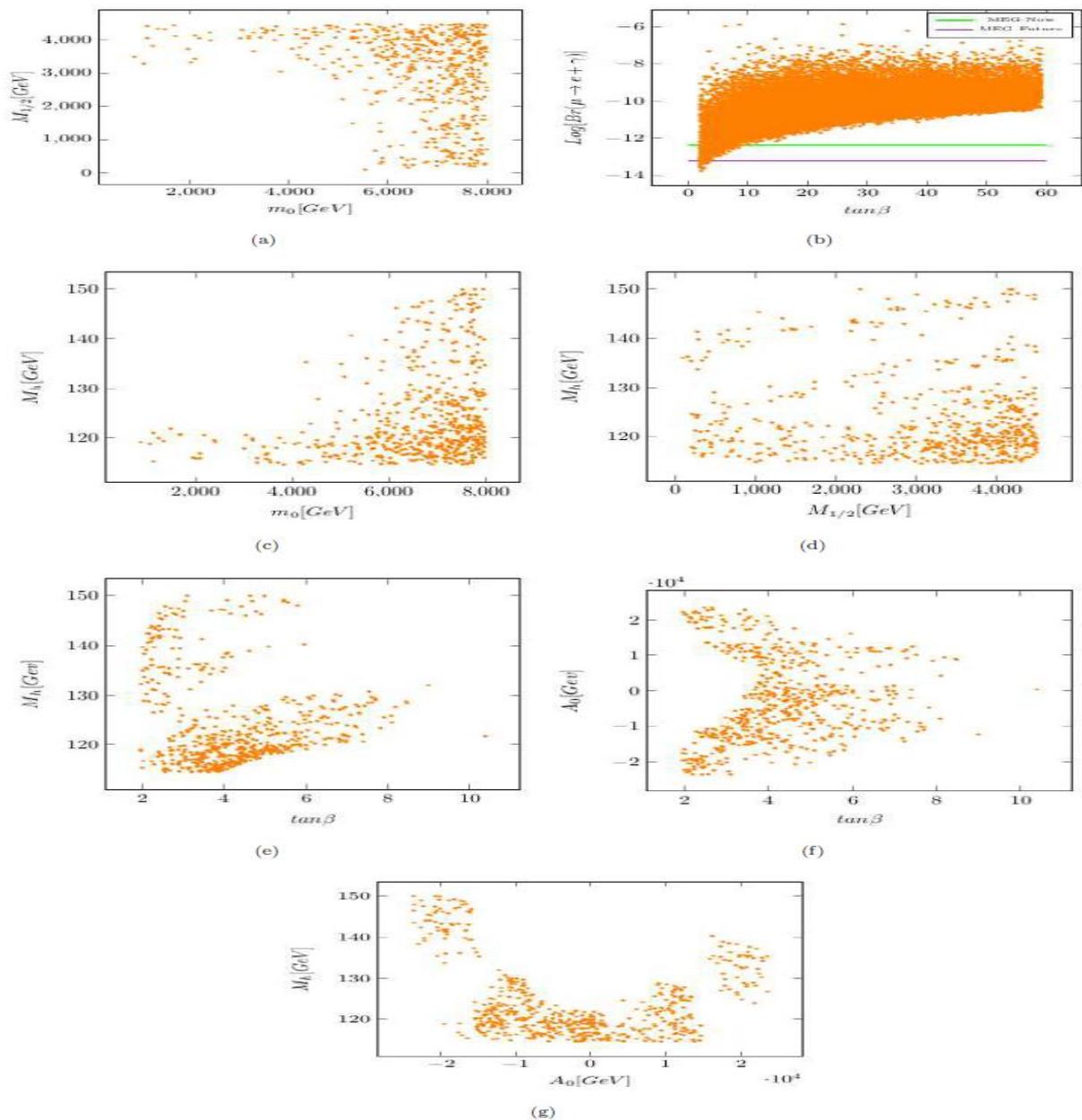


Figure 4: In figs (4a-4e) allowed SUSY parameters region as constrained by MEG 2016 bound is presented.

In figs. 5c, $\text{Log}[BR(\mu \rightarrow e + \gamma)]$ Vs $\tan\beta$ is presented. It is grasped from the fig. 5c, that for present bound of MEG, almost all values of $\tan\beta$ from 5 to 60 are favoured, whereas future MEG bound restricts the values of $\tan\beta$ to 5–40. Fig.5d represents m_0 Vs $M_{1/2}$. The SUSY parameter space $M_{1/2}$ – m_h and m_0 – m_h is presented, as allowed by present MEG bounds in figs. 6a,6c. For Higgs mass to be around 126 GeV, values of $M_{1/2}$ from 4 TeV to 5 TeV are mostly

International Journal of Innovative Research in Science, Engineering and Technology

(A High Impact Factor, Monthly, Peer Reviewed Journal)

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Vol. 9, Issue 3, March 2020

allowed. Similarly for m_h around 126 GeV, region $6 \text{ TeV} \leq m_0 \leq 8 \text{ TeV}$ are mostly allowed. Fig 6b shows A_0 [GeV] V_{smh} [GeV]. Negative values of A_0 are favoured in order to have Higgs mass around 126 GeV. Fig 6d represents $\tan \beta V_{smh}$. The last row in the right panel depicts the constraint restriction on $\tan \beta$. Almost all values of $\tan \beta$ from around 5 to 40 are allowed. In the CMSSM case, the restrictions in the low $\tan \beta$ are due to m_h while those at high $\tan \beta$ are from the constraints on $BR(\mu \rightarrow e + \gamma)$. In the CMSSM case, a light Higgs mass around 125 GeV does not necessarily implies a suppressed flavor violating entry instead of the largeness of A-terms required. In fact flavor violation constraints are indeed still very powerful.

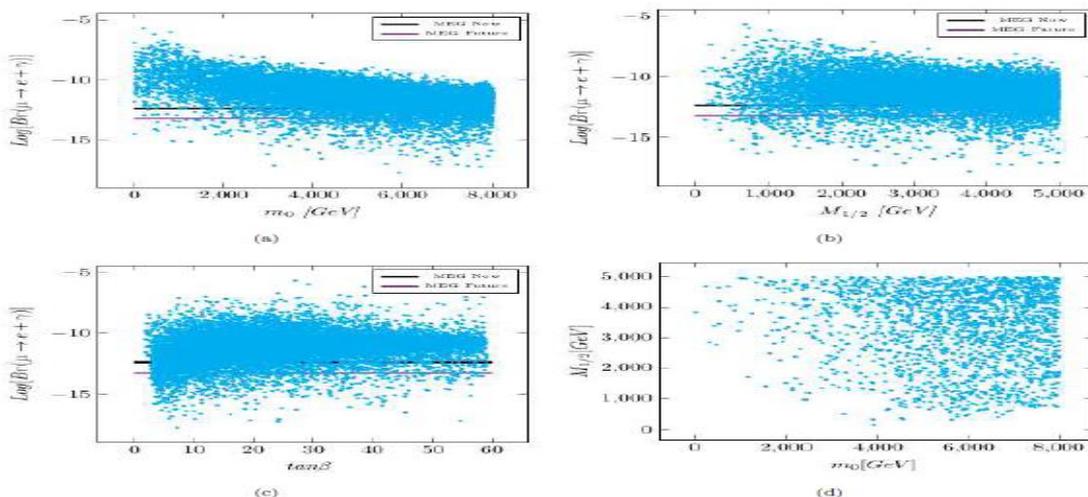


Figure 5: The consequences of the analysis and calculations are presented for NUHM case. In fig 5a,5b, different horizontal lines illustrates the present (MEG 2016) by MEG Collaboration and future MEG bounds for $BR(\mu \rightarrow e\gamma)$. Figs. 5c,5d portrays the allowed SUSY parameter space for different parameters, as is constrained by stringent by MEG 2016 bounds.

C. Non Universal Scalar Mass Models (NUSM)

The parameters of NUSM model is given by [31]

$\tan \beta, M_{1/2}, A_0, \text{sgn}(\mu),$ and m_0 .

The parameters of this model have the exact role to those in CMSSM case except for a major dissimilarity in the scalar sector. First two generations scalars masses (squarks and sleptons) and the third generations sleptons masses are designated as m_0 at the GUT scale. Here m_0 is stretched over a very large value upto tens of TeVs. Nevertheless the Higgs scalars and the third family of squarks are presumed to have dispersed mass values at MGUT. In this computation the mass of third generation of squarks and Higgs scalars are zero. A vanishing A_0 in our analysis is conjectured. [39]. Computations obtained with the non universal scalar masses at M_{GUT} is presented. In fig.7a, the soft SUSY parameter space as allowed by present and future MEG bounds on $BR(\mu \rightarrow e)$ is depicted. For the present MEG bound on $\text{Log}BR(\mu \rightarrow e\gamma)$ i.e -12.376, a large part of $M_{1/2}$ parameters space survives. $M_{1/2}$ lies between $1 \text{ TeV} \leq M_{1/2} \leq 6 \text{ TeV}$. In fig.7b $2 \text{ TeV} \leq m_0 \leq 16 \text{ TeV}$ exists for MEG constraint i.e -12.376. The fig.7c in the right panel shows m_0 GeVs $M_{1/2}$ GeV as constrained by MEG 2016 bound on $BR(\mu \rightarrow e)$. From fig. 7d it is seen that present MEG bound allows $\tan \beta$ to lie between 5 to 47, whereas future MEG bound restricts high values of $\tan \beta$ and limits itself to low values, where $5 \leq \tan \beta \leq 18$. It is found from fig.8a that for Higgs boson mass around 125 GeV, $m_0 \geq 10 \text{ TeV}$ is mostly allowed. For $m_h \approx 125.9 \text{ GeV}$ the parameter space $10000 \text{ GeV} \leq m_0 \leq 16000 \text{ GeV}$ is mostly recommended. From fig. 8b it is observed that for higgs mass range around $125 \pm 2 \text{ GeV}$, $M_{1/2}$ lies between $1.5 \text{ TeV} \leq M_{1/2} \leq 5 \text{ TeV}$. From fig.8c, it is grasped Higgs mass of 125 GeV disfavours values of both low $\tan \beta \leq 15$ and high $\tan \beta \geq 30$. Fig. 8d and fig.8e represents favoured values of softsusy space detectable at future run of HE LHC.

International Journal of Innovative Research in Science, Engineering and Technology

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Vol. 9, Issue 3, March 2020

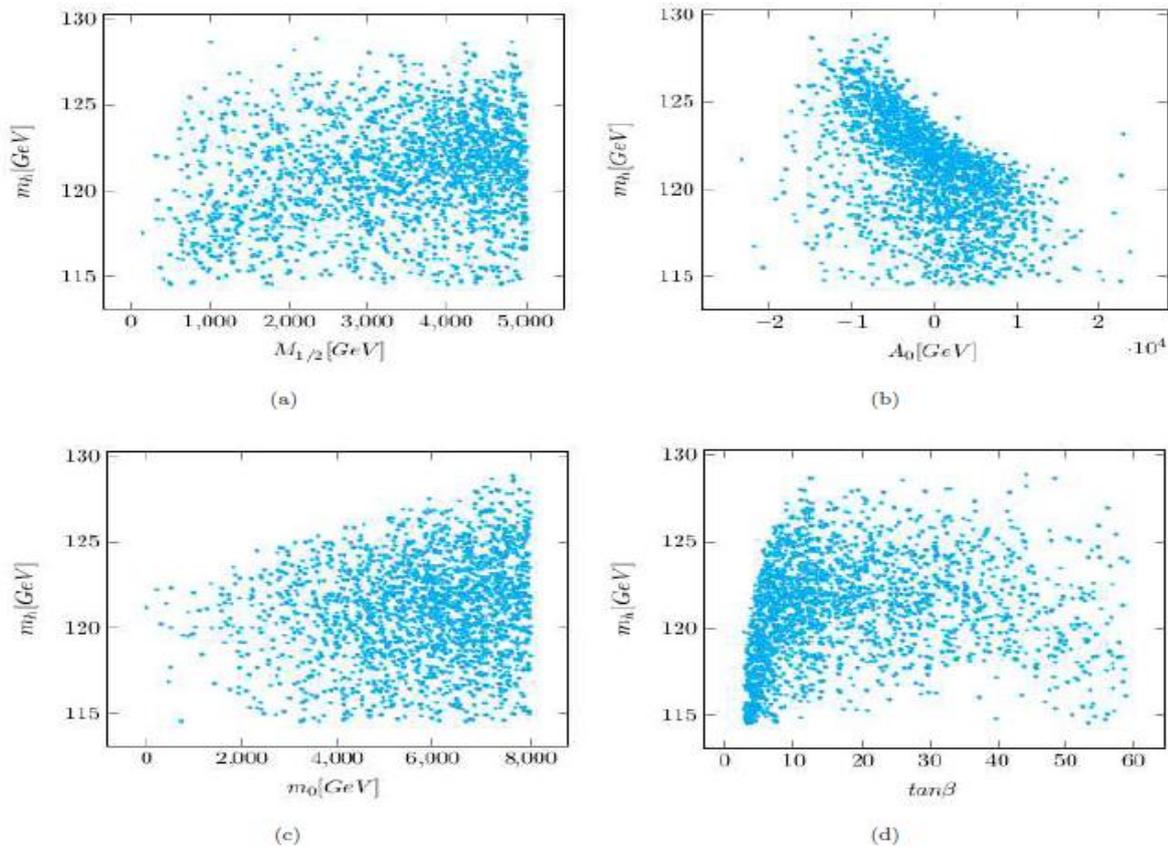


Figure 6: Figs. 6a-6d describe the allowed SUSY region for different soft SUSY parameter space, as is obstructed by stringent MEG 2016 bounds.

In Tables II and III the relative research of the comparative study between different models is presented here. The new consequences in NUSM are the following:

- A. Fragile m_0 SUSY space is allowed as contrasted to CMSSM.
- B. A wider SUSY parameter space is favoured as compared to CMSSM and NUHM model.
- C. For Higgs mass around 125 GeV, favoured values of $\tan \beta$ are $12 \leq \tan \beta \leq 28$ are allowed. $\tan \beta$ values less than 30 are allowed.
- D. The expected sensitivity of the MEG-II experiment which is 6×10^{-14} for three years of data taking restricts values of $\tan \beta$ to be less than 20 and these predict or forecast small LFV rates which is easily accessible from the figures.

International Journal of Innovative Research in Science, Engineering and Technology

(A High Impact Factor, Monthly, Peer Reviewed Journal)

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Vol. 9, Issue 3, March 2020

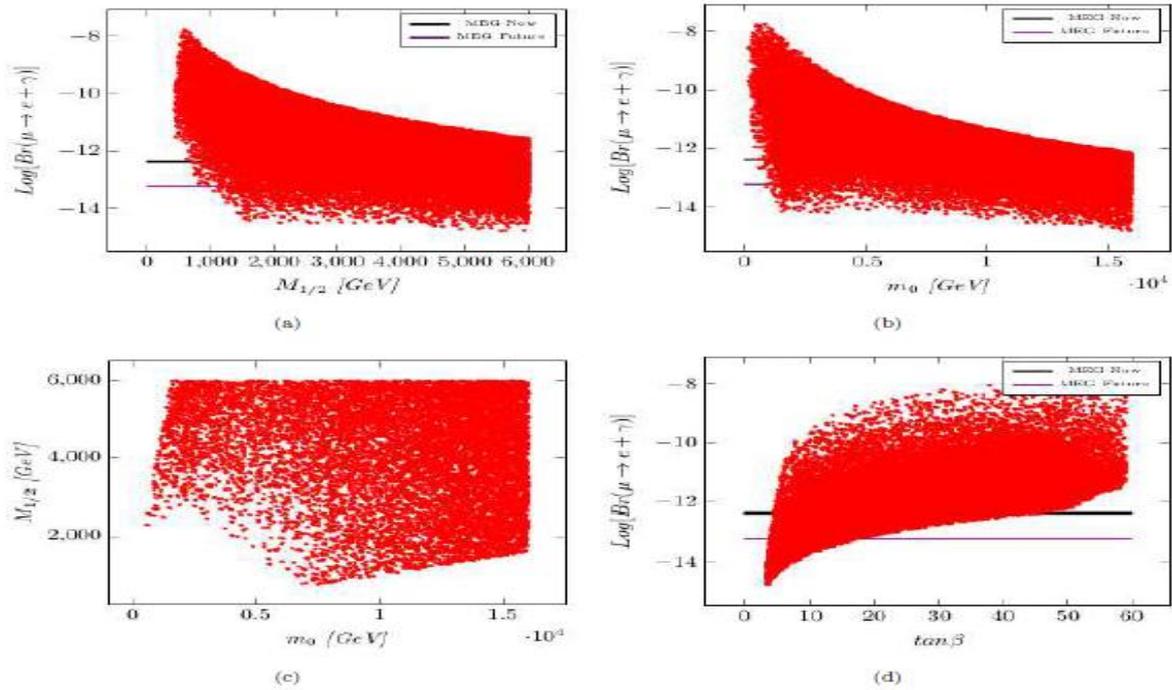


Figure 7: The results of the analysis are presented for NUSM case. In fig 7a, 7b, different horizontal lines depicts the present (MEG 2016) and future MEG bounds for $BR(\mu \rightarrow e + \gamma)$. Figs. 7c, 7d shows the allowed space for different parameters, that is allowed by MEG 2016 bound.

Table II: Sparticle masses in this table are comparison between model CMSSM and NUHM.

Range of parameters allowed by for $BR(\mu \rightarrow e\gamma) < 4.2 \times 10^{-13}$ MEG 2016, for CMSSM	Range of parameters allowed by $BR(\mu \rightarrow e\gamma) < 4.2 \times 10^{-13}$ MEG 2016, for NUHM
1.Fig 3a: $M_{1/2} \geq 3.2$ TeV by MEG 2016 Fig 3b: $m_0 \geq 4.2$ TeV	1.Fig 5b: For MEG 2016, $M_{1/2} \geq 1$ TeV Fig: 5a, $m_0 \geq 2.2$ TeV for $BR(\mu \rightarrow e\gamma) < 10^{-13}$.
2.Fig 4b (MEG 2016), : $\tan \beta$ lies in the range, $2 \leq \tan \beta \leq 10$ for $BR(\mu \rightarrow e\gamma) < 4.2 \times 10^{-13}$	2.Fig. 5c for MEG 2016 of $BR(\mu \rightarrow e\gamma)$ $\tan \beta$ lies in the range, $3 \leq \tan \beta \leq 60$ for $BR(\mu \rightarrow e\gamma) < 4.2 \times 10^{-13}$
3.Fig 4g: A_0 lies in the range, $A_0 \sim -9000$ GeV and $A_0 \sim 12000$ GeV for $m_h = 125.9$ GeV	3. Fig: 6b, A_0 lies in the range, -10000 GeV $\leq A_0 \leq -1000$ GeV, for $m_h = 125.9$ GeV

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(A High Impact Factor, Monthly, Peer Reviewed Journal)

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Vol. 9, Issue 3, March 2020

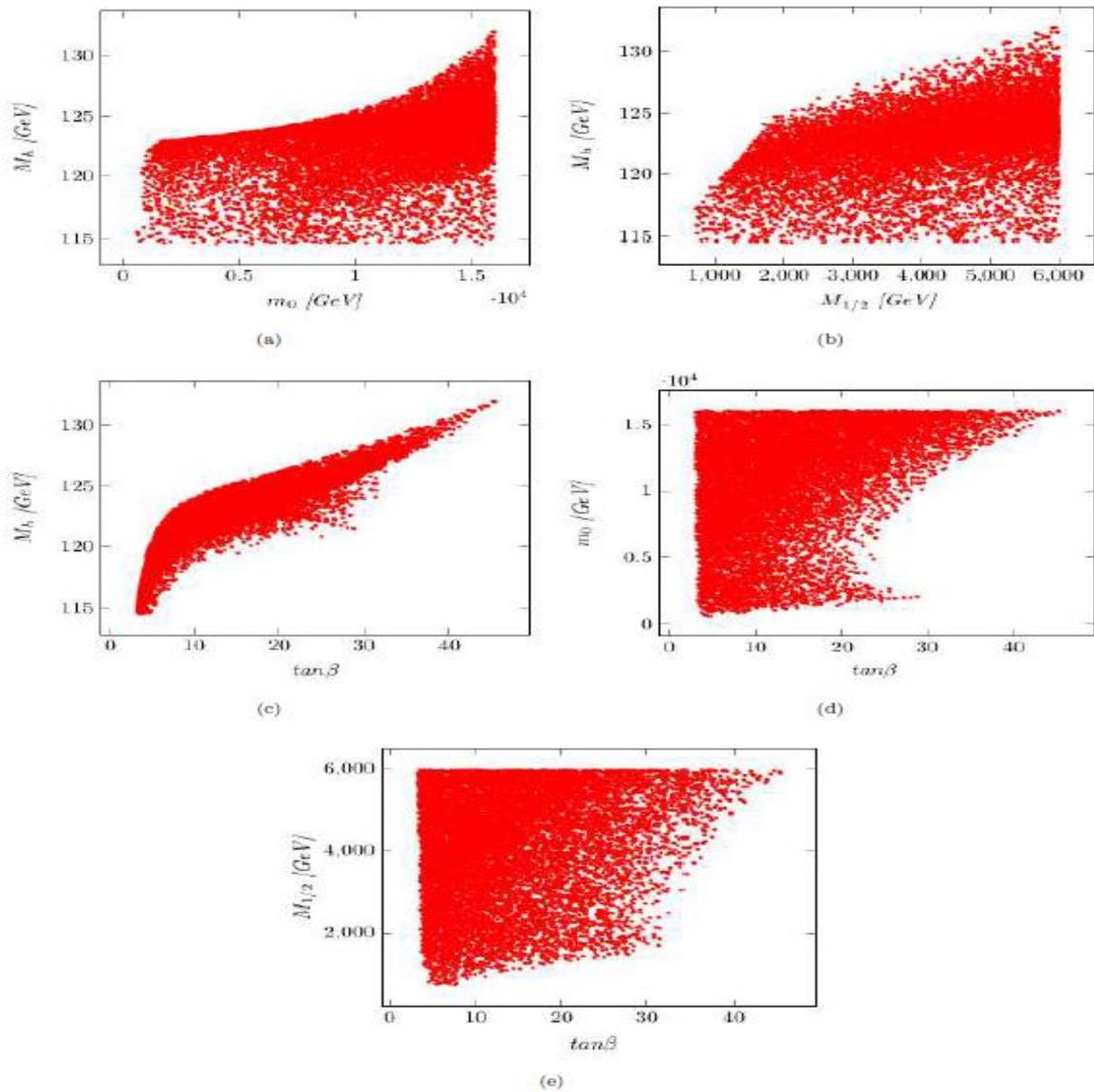


Figure 8: The results of the computations presented for NUSM case. Fig 8a, 8b, 8c, 8d, 8e, depicts the permitted space for different parameters, that is allowed by MEG 2016 bound by MEG collaboration.

International Journal of Innovative Research in Science, Engineering and Technology

(A High Impact Factor, Monthly, Peer Reviewed Journal)

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Vol. 9, Issue 3, March 2020

Table III: Sparticle masses in this table are comparison between model NUSM and NUHM.

Range of parameters allowed by for BR ($\mu \rightarrow e\gamma$) $< 4.2 \times 10^{-13}$ MEG 2016, for CMSSM	Range of parameters allowed by BR ($\mu \rightarrow e\gamma$) $< 4.2 \times 10^{-13}$ MEG 2016, for NUHM
1.Fig 7a: $M_{1/2} \geq 1$ TeV by MEG 2016 Fig 7b: $m_0 \geq 2.2$ TeV	1.Fig 5b: For MEG 2016, $M_{1/2} \geq 1$ TeV Fig: 5a, $m_0 \geq 2.2$ TeV for BR($\mu \rightarrow e\gamma$) $< 10^{-13}$.
2.Fig 7b (MEG 2016), : $\tan \beta$ lies in the range, $5 \leq \tan\beta \leq 50$ for BR ($\mu \rightarrow e\gamma$) $< 4.2 \times 10^{-13}$	2.Fig. 5c for MEG 2016 of BR ($\mu \rightarrow e\gamma$) $\tan \beta$ lies in the range, $3 \leq \tan\beta \leq 60$ for BR ($\mu \rightarrow e\gamma$) $< 4.2 \times 10^{-13}$
3.Fig 8c: $\tan\beta$ lies in the range, $15 \leq \tan\beta \leq 30$ for $m_h = 125.9$ GeV	3. Fig: 6d, $\tan\beta$ lies in the range, $8 \leq \tan\beta \leq 13$ for $m_h = 125.9$ GeV

V. CONCLUSION

The see-saw mechanism is very encouraging in the sense that an intermediate mass scale exists for Majorananeutrinos which solves the mystery of the tiny active neutrino masses nevertheless also describing the absence of righthandedneutrino effects in low energy laboratory scale. The see-saw mechanism is not an impelling extension of theSM since the Higgs boson mass would likely blasts to the see-saw scale owing to its quadratic divergent radiative corrections to the Higgs scalar masses. Supersymmetry stabilizes the hierarchy problem and steadfast the Higgs massso the weak scale can mutually stays along with the Majorana mass scale (and the GUT and Planck scales). Theaim of this paper is to present forecast for LFV process $\mu \rightarrow e + \gamma$ within credible SUSY models that are consistentwith LHC Run 2 results. Within SUSY models, LFV mechanisms should happen, possibly at an observable level.The motivation for this paper is to present predictions for cLFV processes within various SUSY models that arehappening with LHC Run 2 results. These should naturally include SUSY models with radiatively-driven naturalness[22] which requires a 125 GeV Higgs mass along with multi-TeV soft terms (as implied by LHC data) nevertheless atthe same time it avoids the fine-tunings attached with the Little Hierarchy problem. Here, the current and projectedreaches of LFV search $\mu \rightarrow e + \gamma$ in MEG II and MEG collaboration within CMSSM, NUHM and NUSM SUSY modelwhere neutrinos are generated with a type-II seesaw mechanism is examined. The soft SUSY parameter space isconstrained within the HE-LHC reach. The NUHM model is well inspired in that it admits for weak scale naturalnessaccompanying with a 125 GeV Higgs mass and sparticles beyond LHC Run 2 limits. The value of Higgs mass as measured at LHC, latest global data on the reactor mixing angle θ_{13} for neutrinos, andlatest constraints on BR($\mu \rightarrow e$) as projected by MEG collaboration[18],[19] is used here. In CMSSM a very heavy $M_{1/2}$ region is allowed by projected sensitivity of MEG II experiments which is 6×10^{-14} , nevertheless in NUHMcase comparatively a low $M_{1/2}$ region is also favoured. Further the non universal scalar mass model (NUSM) is alstudied. As compared to CMSSM, in NUHM, a broader soft susy space is allowed. It is found that in NUSM, awide spectra of susy parameter space is admitted, as compared to both CMSSM and NUHM. In NUHM, it is seen that favoured values of $|A_0|$ consistent with the newly discovered Higgs boson mass are negative and are shiftedtowards lighter side (compared to CMSSM). It is perceived that in NUSM as Log[BR($\mu \rightarrow e\gamma$)] decreases from presentMEG bound to around to future MEG II projected sensitivity, a wide region of soft susy parameter space of $M_{1/2}$ from 1 TeV to almost 6 TeV is allowed. The computations in this paper show that for natural SUSY with the

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Vol. 9, Issue 3, March 2020

model CMSSM, rates for $BR(\mu \rightarrow e\gamma)$ are already ruled out by the projected MEG II sensitivity experiment for $S_4 \times Z_n$ flavor symmetric SUSY SO(10) theory scenarios. The non universal boundary conditions in NUSM can probe experimental indications for the production of supersymmetric particles and can entertain detector set up to ensure that various super symmetric models can probe signatures at HE LHC. Observation of sparticles at HE/HL LHC, could help us to distinguish among CMSSM, NUSM and NUHM models, in context to the limits put by cLFV decays. This in turn leads to a motivation for examining various SUSY theories beyond standard model.

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